ICFA mini-workshop on Nonlinear dynamics and Collective effects

Phase-merging enhanced harmonic generation and its advances

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Outline

- Introduction, & TGU background
- Phase-merging Enhanced Harmonic Generation (PEHG)
- Advanced concepts
- **TGU** activities at **SINAP**
- **Summary & Outlook**



Transverse Gradient Undulator (TGU)



- □ T. I. Smith, et al., Reducing the sensitivity of a free-electron laser to electron energy. Journal of Applied Physics, 50 (1979) 4580.
- N. Kroll, et al., Theory of the transverse gradient wiggler. IEEE Journal of quantum Electronics, 17 (1981)1496.
- Z. Huang, et al, Compact X-ray FEL from a Laser-Plasma Accelerator Using a Transverse-Gradient Undulator. Physical Review Letter, **109** (2012) 204801.
- G. Fuchert, et al., A novel undulator concept for electron beams with a large energy spread. Nucl. Instr. and Meth. A. **672** (2012) 33.



TGU for FEL resonance compensation

- **TGU** with linear gradient of α
- **D** Sort beam energy by dispersion η

$$\frac{\Delta K}{K_0} = \alpha x \qquad \qquad x = \eta \frac{\Delta \gamma}{\gamma_0}$$

Optimized resonance relationship

$$\eta = \frac{2 + K_0^2}{\alpha K_0^2}$$

- □ 1GeV, 10kA, 1% energy spread;
- 0.1µm emittance; 5 fs (50 pC)
- **5**-m SC undulator , K = 2;
- Transverse gradient $\alpha = 150 \text{ m}^{-1}$
- Radiation wavelength 3.9 nm
- 1cm transverse dispersion

θ (mrad)

Transverse beam size 100×15µm





Z. Huang, et al., PRL, 2012.



Motivation & How the idea starts





FIG. 2 (color online). Optimization of the transverse gradient α of the modulator and the transverse dispersion η of the dogleg by 1D simulation, in order to find the optimal bunching factor of the 30th harmonic for the cooled HGHG.

H. Deng*, C. Feng, PRL (2013).



Phase-merging: single-particle dynamics

We first derive the mechanism behind such kind of schemes from single-particle dynamics. Practically for a given wavelength of the seed laser, the resonant beam energy should be

$$\gamma_r(x) = \gamma_0 + \alpha \eta \frac{K_0^2}{K_0^2 + 2} (\gamma - \gamma_0).$$
 (1)

Consider a resonant and an arbitrary electron (γ_0, θ_0) and (γ, θ_0) at the exit of the TGU modulator, which is the electron (γ_0, θ_0) and $(\gamma, \theta_0 - \Delta \varphi)$ at the entrance of the modulator, respectively. Then,

$$\begin{cases} \gamma_0' = \gamma_0 - \Delta \gamma \sin \theta_0 = \gamma_0 - \Delta \gamma \theta_0 \\ \gamma' = \gamma - \Delta \gamma \sin(\theta_0 - \Delta \varphi/2) = \gamma - \Delta \gamma(\theta_0 - \Delta \varphi/2), \end{cases}$$
(2)

 $\Box \Delta \varphi$ is the phase exchange difference of the arbitrary electron with respect to the resonant one.

$$\Delta \varphi = 4\pi N \frac{(\gamma - \gamma_r)}{\gamma_0},\tag{3}$$

and N represents the period number of the modulator.



Phase-merging: single-particle dynamics

Combining Eq. (2) and Eq. (3), we can easily derive that

$$\frac{\gamma' - \gamma'_0}{\gamma - \gamma_0} = 1 - \frac{2\pi N \Delta \gamma}{\gamma_0} \left(\frac{\alpha \eta K_0^2}{K_0^2 + 2} - 1 \right).$$
(4)

Eq. (4) illustrates a scaling for longitudinal beam phase space control.

- ✓ **Typical HGHG setup**: the local beam energy spread is amplified by a factor of $2\pi N\Delta\gamma/\gamma_0$ which is usually a relatively small number.
- Typical TGU region: when we increase the αη product and make the right hand of Eq. (4) to be unity, the electron beam energy spread is not changed and almost every electron satisfies the FEL resonant condition.
- Phase-merging: if one further increases αη product properly, the right hand of Eq. (4) can be zero. Although it seems that, the electron beam energy spread is suppressed, in fact, all the electrons with the same energy merges to an energy-related longitudinal phase.



Phase-merging: single-particle dynamics

Some practical numbers:

- ✓ Electron beam: E=0.84GeV, 100keV slice energy spread.
- ✓ The modulator parameters: period length 80mm×12, and K=5.8.
- ✓ 265nm seed laser, energy modulation amplitude 500keV.





Phase-merging: Alternative scheme I









C. Feng, H. Deng*, D. Wang, Z. Zhao, New J. Phys. 16 (2014) 043021



Phase-merging: Alternative scheme II



C. Feng, T. Zhang, H. Deng, Z. Zhao*, Phys. Rev. ST-AB. 17 (2014) 070701



Phase-merging Enhanced Harmonic Generation (PEHG)



HGHG bunching $b_n = e^{-\frac{n^2 D^2 \delta^2}{2}} J_n(n D \Delta \gamma)$ **PEHG bunching** $b_n = J_n(nD\Delta\gamma)$

HGHG/EEHG/PEHG bunching

- 1D results: The maximum bunching scales as 0.67/n^{1/3}
- 1D result: The maximum bunching is independent on the energy modulation
- □ A 3D theory and s2e simulation should be done.



PEHG: Zero response to beam energy chirp



G. Wang, C. Feng, H. Deng, T. Zhang, D. Wang*, NIMA, 753 (2014) 56-60.



PEHG: three-dimensional theory

	S	eed	las	er						$\mathbf{A} \mathbf{\nabla}$	I				
			-			5				VÁ					
		$\Lambda \Lambda$								r	noc	dulator	Т	GU	
		$V \Delta$				do	ogle	g							
												4	7		
	1	L_d	0	η	•	[1	L_m	0	0		1	L_T	0	$\tau L_T/2$?]
Rn -	0	1	0	0	Ru -	0	1	0	0	$R_{TCU} \approx$	0	1	0	au	
$\kappa_D =$	0	η	1	ξ_d	$n_M =$	0	0	1	0		τ	$\tau L_T/2$	1	$\tau^2 L_T/6$	6
	0	0	0	1		0	0	h	1		0	0	0	1	
				1	$L + L_T$	+ 2	L_c	$L_T/$	$2h\tau$	$\eta + \tau L$	₀ (1 -	$+ h\xi_d$			
				0	2	!		h	au	au(1	+h	ξ_d)			
				τ	au	L		0)		0				
				0	h_{1}	η		h	ı	1 -	$+ h\xi$	d			

$$b_n = J_n [nAB] e^{-(1/2)n^2 (T_1^2 + T_2^2)}$$
$$T_1 = k_s \tau \sigma_x \quad T_2 = k_s \tau L \sigma_{x'}$$
$$\begin{cases} L = L_d + L_m + L_T/2 \\ L_0 = L_T/2 + L_c \end{cases}$$

$$b_n = J_n [nAB] e^{-n^2 k_s^2 \tau^2 L \varepsilon_x}$$



PEHG: SXFEL start2end results





PEHG: SXFEL start2end results





Two-stage PEHG to hard X-ray ?





PEHG-assisted ultrafast pulse generation

According to beam density modulation theory, current & bunching factor distribution during one seed wavelength can be expressed as



H. Deng et al, Chin. Phys. C 2010 K. Li, C. Feng, H. Deng, et al, 2017, in preparation.



Coherent harmonic generation at storage ring

Main parameters

Beam energy: 600MeV Energy spread: 0.6MeV Emittance: 17.5nm-rad Coupling: 3% Seed laser: 800nm Seed modulation: 1.2MeV Radiation: 133nm



Considering the small vertical emittance, each bunch was proposed to be vertically dispersed only after it undergoes sufficient damping. Then under an optimal condition, the bunching factor of the 6th harmonic is enhanced to 23.0% by PEHG from 1.8% in OK setup.

More advanced schemes are proposed by C. Feng, et al., for EUV lithography



TGU-assisted MBI suppression



C. Feng et al., New J. Phys. 17, 073028 (2015).





D. Huang et al., Phys. Rev. Accel. Beams 19, 100701 (2016).

 T. Liu et al., Phys. Rev.

 Accel.
 Beams
 20,

 082801 (2017).



TGU-assisted MBI suppression





TGU-enhanced transverse deflector



	FLASH	FERMI	LCLS
TDS type	S-band	S-band	X-band
TDS freq.	2856MHz	2998MHz	11424MHz
TDS Voltage	26MV	20MV	48MV
Time resolution	~27fs	~20fs	~1fs



TGU-enhanced transverse deflector



G. Wang, H. Deng*, et al, arXiv:1510.06111.







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TGU-60 modulator prototype (2013)









TGU-20 radiator (2016)



Period length	20	mm
Segment length	1.5	m
Gap	>7.00	mm
Peak field	0.615	Т
К	1.150	
Gradient	50	m⁻¹
Gradient tolerance	±5	m⁻¹





POP experiment of PEHG





Conclusions

- Phase-merging in laser-beam interaction was proposed & studied. Once the door was opened, alternative schemes can be used to achieve the proposed phase-merging phenomenon.
- Phase-merging enhanced FEL is one of the most straightforward applications in seeding business. The 3D analytical theory and s2e simulations were performed, which demonstrates the feasibility of fully coherent soft-x-ray FEL from the commercial laser using singlestage PEHG technique (30th harmonic or even higher).
- Many advanced concepts of phase-exchange, i.e., transverselongitudinal coupling is being studied, i.e., ultra-fast pulse generation, ring-based schemes, enhanced TDS and MBI suppression, etc.
- Several TGUs have been successfully manufactured at SINAP. Some proof-of-principle experiments of PEHG is under consideration, with the funding supports from NSFC and MOST of China.

SXFEL facility



Thanks for attention !